

CHAPTER 10 – ANALYSIS OF MAGNETICALLY COUPLED CIRCUITS

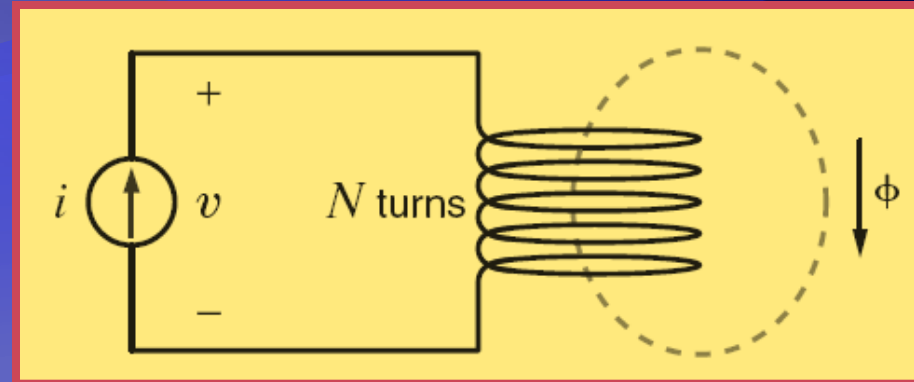
LEARNING GOALS

- 1. To learn how to apply Ampere's law and Faraday's law on magnetically coupled coils**
- 2. To differentiate between self inductance and mutual inductance**
- 3. To learn the relation between mutual inductance and coupling coefficient**
- 4. To calculate voltages and currents in circuits containing mutual inductance**
- 5. To calculate voltages and currents in circuits containing ideal transformers**

INTRODUCTION

Most **kitchen cooktops** used today use **gas or electricity** to heat up **an element**, which then transfers the heat to a pot or pan placed on it. This type of cooktops has an **efficiency of about 40 or 50%** due to the **indirect heat transfer**. Another type of cooktops is the **induction cooktop** which uses an electromagnet excited by a high frequency ac signal. The magnetic field created by the electromagnet induces a **current into the pot or pan and heats it directly**. This type of cooktops has an **efficiency reaching 90%**. In addition, because the pot or pan is heated directly, the **surface of the cooktop remains cool!**

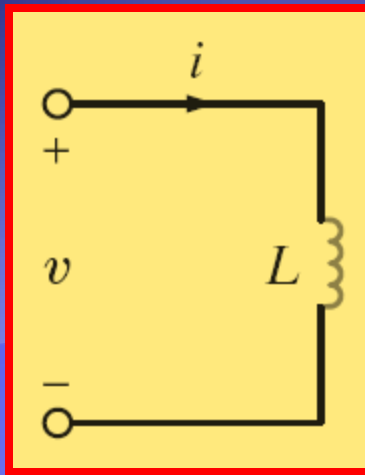
MUTUAL INDUCTANCE



Given a time-varying current, $i(t)$, flowing in an **ideal coil** that has N turns, **Ampere's law** predicts that a **magnetic flux** will be created around the coil, and that the **flux linkage** is proportional to the current that has created it:

$$\lambda(t) = N\phi(t) = Li(t)$$

Where L is the (self) inductance of the coil, and $\phi(t)$ is the magnetic flux.

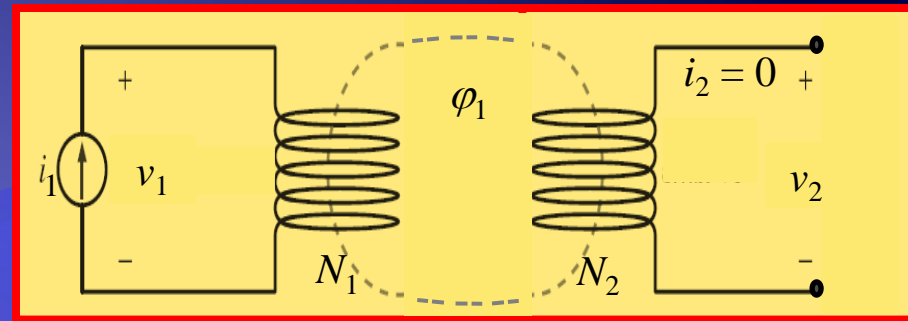


On the other hand, **Faraday's law** predicts that the **flux linkage induces a voltage** across the coil equal to:

$$v(t) = \frac{d\lambda(t)}{dt} = L \frac{di(t)}{dt}$$

where $i(t)$ and $v(t)$ satisfy the **passive sign convention**.

If we now bring a **second ideal, open-ended coil with N_2 turns** close enough to the first coil such that **the flux pro-**



duced by the first coil also links the second coil, we could express the flux linkage of the second coil by:

$$\lambda_2(t) = N_2 \phi_1(t)$$

Where the flux $\phi_1(t)$ is produced by the current $i_1(t)$ flowing in the first coil.

$$\lambda_1(t) = N_1 \phi_1(t) = L_1 i_1(t)$$

Consequently, **although no current is flowing in the second coil**, a voltage across it will be induced:

$$v_2(t) = \frac{d\lambda_2(t)}{dt} = N_2 \frac{d\phi_1(t)}{dt}$$

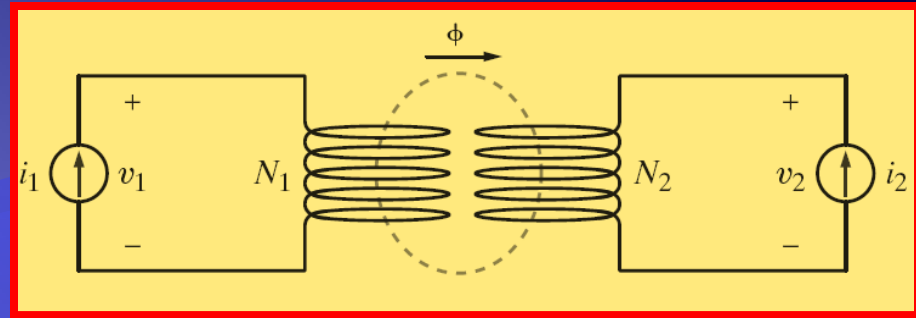
Using the expression of $\varphi_1(t)$ produced by $i_1(t)$:

$$\varphi_1(t) = \frac{L_1}{N_1} i_1(t)$$

we deduce the expression of **the voltage $v_2(t)$ induced by current $i_1(t)$ across the second coil**:

$$v_2(t) = \frac{N_2}{N_1} L_1 \frac{di_1(t)}{dt} \equiv L_{21} \frac{di_1(t)}{dt}$$

As we can see, **the induced voltage across the second coil is proportional to rate of change of the current flowing in the first coil**. The constant of proportionality is called the **mutual inductance** and is expressed in **Henrys**. We say that the two coils are **magnetically coupled**.



At this stage, **if we connect a current source to the terminals of the second coil**, it will not only create a voltage across the second coil, but **it will also induce an additional voltage across the first coil**, due to the **magnetic coupling** between the two coils.

In the case where both currents, **$i_1(t)$ and $i_2(t)$, contribute to the magnetic flux, $\varphi(t)$** (which depends on the respective orientations of the coil windings), the flux linkages for each coil could be expressed as:

$$\lambda_1 = L_1 i_1 + L_{12} i_2$$

$$\lambda_2 = L_2 i_2 + L_{21} i_1$$

And the resulting expressions of the voltages across each coil become:

$$v_1 = \frac{d\lambda_1}{dt} = L_1 \frac{di_1}{dt} + L_{12} \frac{di_2}{dt}$$

$$v_2 = \frac{d\lambda_2}{dt} = L_2 \frac{di_2}{dt} + L_{21} \frac{di_1}{dt}$$

For **linear circuits**, we generally consider that:

$$L_{12} = L_{21} = M$$

Where M is the **mutual inductance**.

We can see that the voltage across each coil is now composed of two terms:

- A **self-induced term** of the form $L_i \frac{di_i}{dt}$
- A **mutually-induced term** of the form:

$$M \frac{di_j}{dt}$$

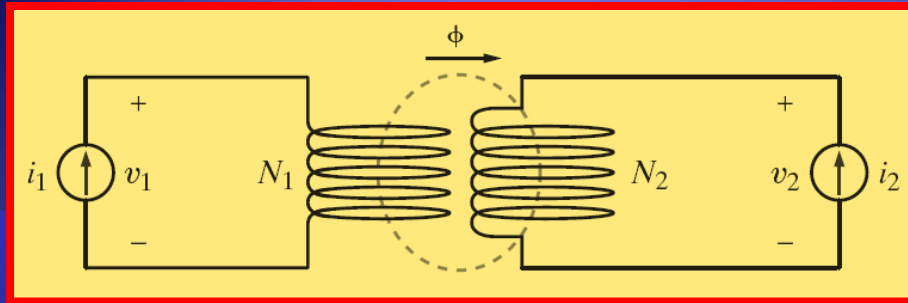
where i and j are 1 and 2.

We note that **if the current $i_2(t)$ is reversed**, its induced flux will oppose the magnetic flux induced by $i_1(t)$. In this case, the expressions of the voltages across each coil become:

$$v_1 = L_1 \frac{di_1}{dt} - M \frac{di_2}{dt}$$

$$v_2 = L_2 \frac{di_2}{dt} - M \frac{di_1}{dt}$$

THE DOT CONVENTION



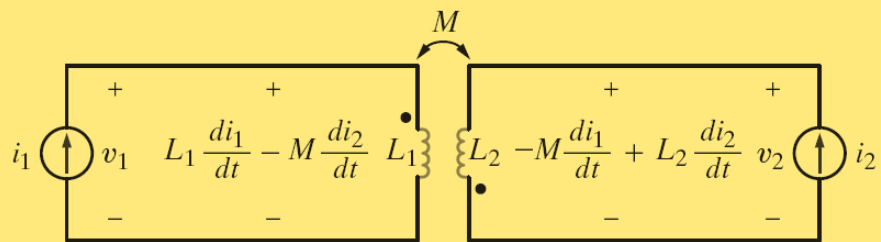
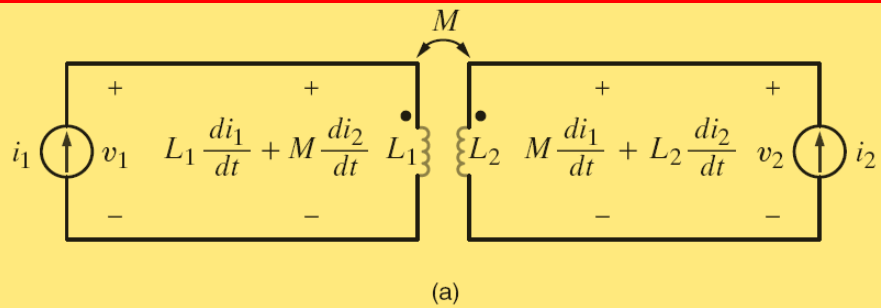
We note that a **negative mutually-induced voltage** will also result if the **orientation of the winding** of one of the coils is **reversed**, as shown in the Figure.

Since the representation of the winding orientation of coils in circuit diagrams may be complicated, we adopt the **dot convention** to take into consideration the **winding orientation** of the coils.

To each coil we associate a dot placed on one of its terminals such that **if the current flows into the dot** of this coil it induces a **mutual voltage** into the second coil that is **positive at the dot of the second coil**. Conversely, **if the current in the first coil flows out of its dot**, it induces into the second coil a **voltage that is negative at the dot of the second coil**.

It is to be noted that:

- the **self-induced voltage follows the passive sign convention**
- while the **mutually-induced voltage follows the dot convention**.



In the first circuit diagram above, **both currents flow into the dots** of their respective coils. Therefore, the circuit diagram is a representation of the case where the current $i_1(t)$ **contributes to the flux produced by $i_2(t)$** , and vice-versa, i.e. the mutually-induced voltage produced by $i_1(t)$ across coil L_2 is positive at the dot of coil L_2 and vice-versa.

However, in the second circuit diagram the current in coil L_1 flows into its dot while the **current in coil L_2 flows out of its dot**. The circuit diagram is therefore a representation of the case where the current $i_1(t)$ **produces a flux that is opposed to the flux produced by $i_2(t)$** , and vice versa, i.e. the mutually-induced voltage produced by $i_1(t)$ across coil L_2 is negative at the dot of coil L_2 and vice-versa.

PROBLEM SOLVING STRATEGY

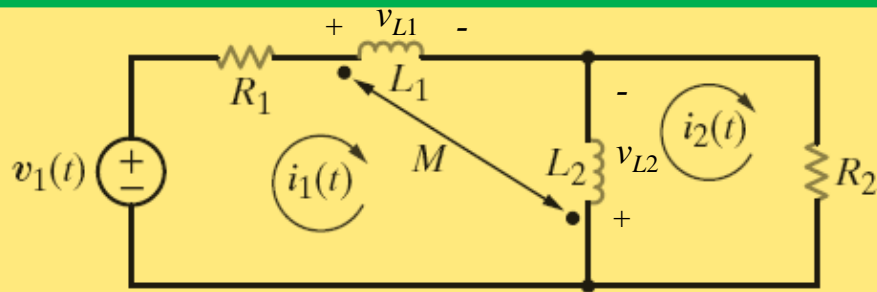
Step 1: Assign **mesh currents** and **coil voltages** and write **coupling equations** using dot convention

Step 2: Write **mesh equations** using mesh currents and coil voltages

Step 3: **Solve** for currents and voltages.

Example 10.1 – Determination of coupling equations

Example 10.2 – Mesh equations for magnetically coupled and electrically connected coils



Mesh equations:

$$-v_1(t) + R_1 i_1(t) + v_{L1}(t) - v_{L2}(t) = 0$$

$$v_{L2}(t) + R_2 i_2(t) = 0$$

Coupling equations:

$$v_{L1}(t) = L_1 \frac{di_1(t)}{dt} + M \frac{d(i_2(t) - i_1(t))}{dt}$$

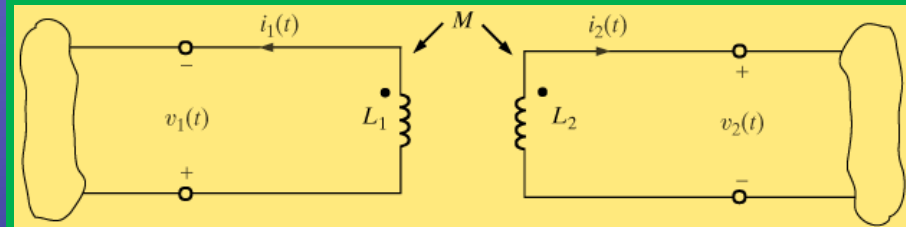
$$v_{L2}(t) = L_2 \frac{d(i_2(t) - i_1(t))}{dt} + M \frac{di_1(t)}{dt}$$

Resulting equations:

$$v_1(t) = R_1 i_1(t) + L_1 \frac{di_1(t)}{dt} + M \frac{d(i_2(t) - i_1(t))}{dt} - L_2 \frac{d(i_2(t) - i_1(t))}{dt} - M \frac{di_1(t)}{dt}$$

$$L_2 \frac{d(i_2(t) - i_1(t))}{dt} + M \frac{di_1(t)}{dt} + R_2 i_2(t) = 0$$

Exercise E10.1 – Coupling Equations for magnetically coupled coils



$$v_1(t) = L_1 \frac{di_1(t)}{dt} + M \frac{di_2(t)}{dt}$$

$$v_2(t) = -L_2 \frac{di_2(t)}{dt} - M \frac{di_1(t)}{dt}$$